1. Convert the covariant expression  $\bar{\Psi}\rho^{\alpha}\partial_{\alpha}\Psi$ , where

$$\psi = \begin{pmatrix} \psi_+ \\ \psi_- \end{pmatrix} \tag{1}$$

and  $\bar{\Psi} = \Psi^{\dagger} \rho^{0}$  are two-dimensional spinors, into the component expression

$$-2(\Psi_+\partial_-\Psi_+ + \Psi_-\partial_+\Psi_-) \tag{2}$$

with  $\partial_{\pm} = \frac{1}{2}(\partial_0 \pm \partial_1).$ 

Hint: Use the following convenient basis for the Dirac matrices

$$\rho^0 = \begin{pmatrix} 0 & 1\\ -1 & 0 \end{pmatrix}, \qquad \rho^1 = \begin{pmatrix} 0 & 1\\ 1 & 0 \end{pmatrix}$$
(3)

2. Prove the following property (valid in two dimensions)

$$\rho^{\alpha}\rho_{\beta}\rho_{\alpha} = 0 \tag{4}$$

3. Show that the action

$$S = -\frac{1}{8\pi} \int d^2 \sigma e \left[ \frac{2}{\alpha'} h^{\alpha\beta} \partial_\alpha X^\mu \partial_\beta X_\mu + 2i \bar{\psi}^\mu \rho^\alpha \partial_\alpha \psi_\mu - i \bar{\chi}_\alpha \rho^\beta \rho^\alpha \psi^\mu \left( \sqrt{\frac{2}{\alpha'}} \partial_\beta X_\mu - \frac{i}{4} \bar{\chi}_\beta \psi_\mu \right) \right]$$
(5)

is invariant under the following local world-sheet symmetries:

(a) Supersymmetry

$$\sqrt{\frac{2}{\alpha'}} \delta_{\epsilon} X^{\mu} = i \bar{\epsilon} \psi^{\mu},$$

$$\delta_{\epsilon} \psi^{\mu} = \frac{1}{2} \rho^{\alpha} \left( \sqrt{\frac{2}{\alpha'}} \psi^{\mu} \partial_{\alpha} - \frac{i}{2} \bar{\chi}_{\alpha} \right) \epsilon,$$

$$\delta_{\epsilon} e_{\alpha}{}^{a} = \frac{i}{2} \bar{\epsilon} \rho^{\alpha} \chi_{\alpha},$$

$$\delta_{\epsilon} \chi_{\alpha} = 2D_{\alpha} \epsilon$$

where  $\epsilon(\tau, \sigma)$  is a Majorana spinor which parametrizes susy transformations and  $D_{\alpha}$  is a covariant derivative with torsion:

$$D_{\alpha}\epsilon = \partial_{\alpha}\epsilon - \frac{1}{2}\omega_{\alpha}\bar{\rho}\epsilon$$
$$\omega_{\alpha} = -\frac{1}{2}\epsilon^{ab}\omega_{\alpha ab} = \omega_{\alpha}(e) + \frac{i}{4}\bar{\chi}_{\alpha}\bar{\rho}\rho^{\beta}\chi_{\beta}$$
$$\omega_{\alpha}(e) = -\frac{1}{e}e_{\alpha a}\epsilon^{\beta\gamma}\partial_{\beta}e_{\gamma}{}^{a}$$

where  $\omega_{\alpha}(e)$  is the spin connection without torsion.

(b) Weyl transformations

$$\begin{split} \delta_{\Lambda} X^{\mu} &= 0 \\ \delta_{\Lambda} \psi^{\mu} &= -\frac{1}{2} \Lambda \psi^{\mu} \\ \delta_{\Lambda} e_{\alpha}{}^{a} &= \Lambda e_{\alpha}{}^{a} \\ \delta_{\Lambda} \chi_{\alpha} &= \frac{1}{2} \Lambda \chi_{\alpha} \end{split}$$

(c) Super-Weyl transformations

$$\delta_{\eta}\chi_{\alpha} = \rho_{\alpha}\eta$$
  
$$\delta_{\eta}(\text{others}) = 0$$

with  $\eta(\tau, \sigma)$  a Majorana spinor parameter.

(d) Two-dimensional Lorentz transformations

$$\begin{split} \delta_l X^{\mu} &= 0\\ \delta_l \psi^{\mu} &= -\frac{1}{2} l \bar{\rho} \psi^{\mu}\\ \delta_l e_{\alpha}{}^a &= l \epsilon^a{}_b e_{\alpha}{}^b\\ \delta_l \chi_{\alpha} &= -\frac{1}{2} l \bar{\rho} \chi_{\alpha} \end{split}$$

(e) Reparametrizations

$$\begin{aligned} \delta_{\xi} X^{\mu} &= -\xi^{\beta} \partial_{\beta} X^{\mu} \\ \delta_{\xi} \psi^{\mu} &= -\xi^{\beta} \partial_{\beta} \psi^{\mu} \\ \delta_{\xi} e_{\alpha}{}^{a} &= -\xi^{\beta} \partial_{\beta} e_{\alpha}{}^{a} - e_{\beta}{}^{a} \partial_{\alpha} \xi^{\beta} \\ \delta_{\xi} \chi_{\alpha} &= -\xi^{\beta} \partial_{\beta} \chi_{\alpha} - \chi_{\beta} \partial_{\alpha} \xi^{\beta} \end{aligned}$$

Note that  $\Lambda, \xi$  and l are infinitesimal functions of  $(\tau, \sigma)$ .

(f) Global space-time Poincaré transformations:

$$\delta X^{\mu} = a^{\mu}{}_{\nu}X^{\nu} + b^{\mu}, \qquad a_{\mu\nu} = -a_{\nu\mu}$$
  
$$\delta h_{\alpha\beta} = 0$$
  
$$\delta \psi^{\mu} = a^{\mu}{}_{\nu}\psi^{\nu}$$
  
$$\delta \chi_{\alpha} = 0$$

- 4. Show that the commutator of two supersymmetry transformations on the fields  $X^{\mu}$  and  $\psi^{\mu}$ , i.e.  $[\delta_1, \delta_2] X^{\mu}$  and  $[\delta_1, \delta_2] \psi^{\mu}$ , gives a translation.
- 5. Find the equations of motion for the fields derived from the action (5). Show the tracelessness of the energy momentum tensor  $T_{\alpha\beta}$  and the  $\rho$ -tracelessness of its superpartner  $T_F$ .

- 6. Write the non-vanishing components of the energy-momentum tensor and the supercurrent in light-cone coordinates and show that they are conserved.
- 7. Show that the NS and R boundary conditions on the fermions allow to cancel the boundary terms in the variation of the action.
- 8. Obtain the mode expansions for the fermionic fields in the Ramond and Neveu-Schwarz sectors of the open string.
- 9. Obtain the superVirasoro generators  $L_m$  and  $G_r$  in terms of oscillators.
- 10. Obtain the classical superVirasoro algebra:

$$[L_m, L_n] = -i(m-n)L_{m+n} [L_m, G_r] = -i\left(\frac{1}{2}m - r\right)G_{m+r} \{G_r, G_s\} = -2iL_{r+s}$$
 (6)

Useful identities:

1.  $\bar{\lambda}_1 \rho^{\alpha_1} \cdots \rho^{\alpha_n} \lambda_2 = (-1)^n \bar{\lambda}_2 \rho^{\alpha_n} \cdots \rho^{\alpha_1} \lambda_1$ 2.  $(\bar{\psi}\lambda)(\bar{\phi}\chi) = -\frac{1}{2} \{ (\bar{\psi}\chi)(\bar{\phi}\lambda) + (\bar{\psi}\bar{\rho}\chi)(\bar{\phi}\bar{\rho}\bar{\lambda}) + (\bar{\psi}\rho^{\alpha}\chi)(\bar{\phi}\rho_{\alpha}\lambda) \}$