

# Ergodicity of infinite particle systems and Applications

(Based upon joint works with Zegarlinski, Inglis, Futorny, Bock)

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# Outline

Conservative Interacting Particle Systems

Jordan-Schwinger map for Lie algebras

J.-S. map for topological algebras

# Conservative interacting particle systems (CIPS): Preliminaries

- ▶ Space of configurations  $X$ . Examples:  $\mathbb{R}^{\mathbb{Z}^d}$ ,  $\{0, 1\}^{\mathbb{Z}^d}$  ...
- ▶ Energy function  $E : X \rightarrow \mathbb{R}$ .
- ▶ Equilibrium is described by the family of measures  
 $d\mu_T = " \frac{e^{-\frac{E}{T}}}{N(T)} dx "$  enumerated by the temperature  $T \in A$ .
- ▶ Evolution is described by the Markov semigroup  
 $P_t : L^2(X, d\mu_T) \rightarrow L^2(X, d\mu_T)$ ,  $t \geq 0$  with infinitesimal generator

$$\mathcal{L} = \sum_{\gamma} D_{\gamma}^2 + Z$$

and invariant measures  $\mu_T$ ,  $T \in A$  ( $\mathcal{L}E = 0$ ).

Question of interest: Rigorous proof that measures are invariant? Ergodicity of the system? Rate of Convergence to equilibrium?

## Heat conduction model

N, Zegarlinski, Inglis have considered

$$\mathcal{L} = b^2 \sum_{\mathbf{i} \in \mathbb{Z}^N} \sum_{\mathbf{j} \in \mathbb{Z}^N : |\mathbf{i} - \mathbf{j}|_1 = 1} (x_{\mathbf{i}} \partial_{\mathbf{j}} - x_{\mathbf{j}} \partial_{\mathbf{i}})^2, \quad (1)$$

-infinitesimal generator corresponding to SDE ( $N=1$ )

$$dx_k(t) = -x_k dt + x_{k+1} dW_k(t) - x_{k-1} dW_{k-1}(t), k \in \mathbb{Z} \quad (2)$$

Physical interpretation (Giardina, Kurchan, Redig (2007)):

We have set of atoms on the lattice  $\mathbb{Z}$  described by their momentums  $(x_i)_{i \in \mathbb{Z}}$ ,  $x_i^2$  – energy of the atom  $i$

Operators  $\mathbf{X}_{i,j} = (x_{\mathbf{i}} \partial_{\mathbf{j}} - x_{\mathbf{j}} \partial_{\mathbf{i}})$  conserve local energy  $x_{\mathbf{i}}^2 + x_{\mathbf{j}}^2$ ;

Total energy  $E = \sum_i x_i^2$  is conserved.

Operators  $(x_{\mathbf{i}} \partial_{\mathbf{j}} - x_{\mathbf{j}} \partial_{\mathbf{i}})^2$  describe the exchange of momentum between atoms  $x_{\mathbf{i}}$  and  $x_{\mathbf{j}}$ .

Invariant measures of the form " $d\mu_T = \frac{1}{N_T!} e^{-\frac{E}{T}} dx$ " i.e. gaussian product measures parameterized by  $T$  – "temperature".

$\mathbf{X}_{i,j}$  – nontrivial Lie bracket, Hörmander condition is violated.

# Existence and non-uniqueness of invariant measure

## Theorem

Measures  $\delta_0$  and  $d\mu_T$ ,  $T > 0$  are invariant measures for semigroup  $P_t$ ,  $t \geq 0$ . Furthermore,

$$\int_E f P_t g d\mu_T = \int_E g P_t f d\mu_T, t, T > 0, f, g \in L^2(d\mu_T). \quad (3)$$



## Sketch of the proof.

Trotter formula and representation of generator as a sum of two operators with "simple" structure. □



# Ergodicity with polynomial decay. Part I

Let  $\mathbb{X}$  – Sobolev space of order 1 with homogeneous norm

$$\mathcal{A}(f) \equiv \left( \sum_{i \in \mathbb{Z}^N} \mu_T |\partial_i f|^2 \right)^{1/2} \sim |f|_{\tilde{\mathbb{X}}}.$$

**Lemma (N, Zegarlinski, Inglis)**

For any  $f \in \mathbb{X}$ ,  $i \in \mathbb{Z}^N$  and  $t > 0$ ,

$$\mu_T(|\partial_i(P_t f)|^2) \leq \frac{A^N}{t^{\frac{N}{2}}} \mathcal{A}^2(f), \quad (4)$$

where  $A = \frac{1}{b} \sup_{t>0} \sqrt{t} \int_0^1 e^{-2t(1-\cos(2\pi\beta))} d\beta$ .

The order of convergence above is optimal ( $f = x_i^2$ ). If

$f_K = \sum_{|i| \leq K} x_i^2$  then

$$\mu_T(P_t f_K - \mu_T f_K)^2 \sim C_2(\delta) \frac{1}{\sqrt{t}^N} \mu_T(f_K - \mu_T f_K)^2 + C_3(\mu_T, \delta) e^{-4\delta t}$$

## Proof of Lemma

Denote  $f_t = P_t f$  for  $t \geq 0$ .

$$\begin{aligned} |\partial_i f_t|^2 - P_t |\partial_i f|^2 &= \int_0^t \frac{d}{ds} P_{t-s} |\partial_i f_s|^2 ds \\ &= \int_0^t P_{t-s} (-\mathcal{L}(|\partial_i f_s|^2) + 2\partial_i f_s \mathcal{L} \partial_i f_s + 2\partial_i f_s [\partial_i, \mathcal{L}] f_s) ds \end{aligned}$$

Algebraic structure—

$$[\partial_i, \mathbf{X}_{k,l}] = \delta_{i,k} \partial_k - \delta_{i,l} \partial_l$$

Question: what is this structure?

Applying it we can calculate  $[\partial_i, \mathcal{L}]$  and (after some cancellations) we get

$$F(t) \leq F(0) + \int_0^t b^2 \Delta F(s) ds, \quad t \in [0, \infty).$$

where  $F(\mathbf{i}, t) = \mu_T |\partial_i(P_t f)|^2$  for  $t \geq 0, \mathbf{i} \in \mathbb{Z}^N$ .

$$F(t) \leq e^{tb^2 \Delta} F(0), t \in [0, \infty)$$

## Ergodicity with polynomial decay. Part II

$$\mathcal{A}(f) \equiv \left( \sum_{\mathbf{i} \in \mathbb{Z}^N} \mu_T |\partial_{\mathbf{i}} f|^2 \right)^{1/2}, \quad \mathcal{B}(f) \equiv \left( \sum_{\mathbf{i} \in \mathbb{Z}^N} (\mu_T |\partial_{\mathbf{i}} f|^2)^{\frac{1}{2}} \right).$$

**Corollary (N, Zegarlinski, Inglis)**

For  $f \in \mathbb{X}$  such that  $\mathcal{B}(f) < \infty$ , we have

$$\sum_{\mathbf{i} \in \mathbb{Z}^N} \mu_T |\partial_{\mathbf{i}} (P_t f)|^2 \leq \frac{A^{\frac{N}{2}}}{t^{\frac{N}{4}}} \mathcal{A}(f) \mathcal{B}(f). \quad (5)$$

Furthermore, there exists a constant  $C \in (0, \infty)$  such that

$$\mu_T \left( (P_t f)^2 \log \frac{(P_t f)^2}{\mu_T (P_t f)^2} \right) \leq C \frac{A^{\frac{N}{2}}}{t^{\frac{N}{4}}} \mathcal{A}(f) \mathcal{B}(f), \quad (6)$$

and hence

$$\mu_T (P_t f - \mu_T(f))^2 \leq C \frac{A^{\frac{N}{2}}}{t^{\frac{N}{4}}} \mathcal{A}(f) \mathcal{B}(f), \quad (7)$$

i. e. our system is ergodic with polynomial rate of convergence.

## Algebraic structure



$$[\partial_{\mathbf{i}}, \mathbf{X}_{k,l}] = \delta_{\mathbf{i},\mathbf{k}} \partial_{\mathbf{k}} - \delta_{\mathbf{i},\mathbf{l}} \partial_{\mathbf{l}}$$

The Jordan–Schwinger map  $D$  for a matrix  $X = (X_{ij})_{i,j=1}^n$  (pp. 212–213, Bidernharn-Louck-1981):

$$D : X \mapsto \sum_{i,j=1}^n X_{ij} a_i a_j^*, \quad (8)$$

$\{a_i, a_j^*\}_{i,j=1}^n$  – boson creation and annihilation operators.

Realization:  $a_i = x_i$ ,  $a_j^* = \partial_{x_j}$ . Properties:

- ▶ algebraic homomorphism i.e.  $D([A, B]) = [D(A), D(B)]$ .
- ▶ corresponds to linear vector field.

Physicists use to construct representations of classical Lie algebras of matrices as an operators on Fock space.

# Generalization of Jordan-Schwinger map

## Definition

$V$  and  $W$  – topological vector spaces(TVS) in duality  
 $\langle \cdot, \cdot \rangle_{V,W}$ ,  $V$  is separable, biorthogonal system  $\{e_k, f_k\}_{k=1}^\infty$

$$D(A) := \sum_{\alpha, \beta \in \mathbb{N}} \langle Ae_\alpha, f_\beta \rangle x_\alpha \frac{\partial}{\partial x_\beta}, A \in \mathcal{L}(V, V),$$

$$\partial(h) := \sum_{\alpha \in \mathbb{N}} \langle h, f_\alpha \rangle \frac{\partial}{\partial x_\alpha}, h \in V.$$

$$\bar{\partial}(r) := \sum_{\alpha \in \mathbb{N}} \langle e_\alpha, r \rangle x_\alpha, r \in W.$$

## Lemma

$$[D(A), D(B)] = D([B, A]),$$

$$[\partial(f), D(A)] = \partial(Af),$$

$$[\partial(f), \partial(g)] = 0$$

$$[D(A), \bar{\partial}(r)] = \bar{\partial}(A^*r), f, g \in V, r \in W, A, B \in \mathcal{L}(V, V).$$

Notice that  $X_{i,j} = x_i \partial_{x_j} - x_j \partial_{x_i} = D_{ij} = D(E_{ij} - E_{ji})$  ( $E_{ij}$  –matrix unit),  $\partial_{x_i} = \partial(e_i)$ . So our algebraic structure is an example of identity (9).

# Representation of Lie algebra with linear vector fields

## Theorem

Let  $\mathfrak{g}$  be an arbitrary Lie algebra and  $\rho : \mathfrak{g} \rightarrow \text{End}(V)$  a faithful representation of  $\mathfrak{g}$ . Then  $-D \circ \rho$  and  $D^* \circ \rho$  embed  $\mathfrak{g}$  into  $\widehat{\mathcal{A}}_I$ , and hence, define representations of  $\mathfrak{g}$  by linear vector fields.

$\mathfrak{g}$ -Lie algebra with center  $\text{Cent}(\mathfrak{g})$   $V = \mathfrak{g}$  and restrict the mapping  $D = D_V$  on the subspace

$$Z := \{ad(v) := [v, \cdot], v \in V\} \subset \mathcal{L}(V, V)$$

$$\tilde{D} := D \circ ad : \mathfrak{g} \rightarrow \text{End}(\mathbb{R}[\bar{x}]), \tilde{D}(v) = \sum_{\alpha, \beta \in A} \langle [v, e_\alpha], f_\beta \rangle x_\alpha \frac{\partial}{\partial x_\beta} \quad (9)$$

## Theorem

For any topological locally convex Lie algebra  $\mathfrak{g}$  there exists an embedding, given by formula (9), of  $\mathfrak{g}/\text{Cent}(\mathfrak{g})$  into the semidirect product  $\tilde{D}(\mathfrak{g}) \ltimes \partial(\mathfrak{g}) \subset \text{End}(\mathbb{R}[\bar{x}])$  of linear differential operators.

# Schrödinger-Virasoro Lie algebra

Schrödinger-Virasoro Lie algebras— stat. physics  
 (Henkel-1994).

$s = 0, \frac{1}{2}, \rho \in \mathbb{Q}$ .  $\mathcal{L}[s, \rho]$ —Lie algebra  $\{L_n, Y_p, n \in \mathbb{Z}, p \in \mathbb{Z} + s\}$

$$\begin{aligned} [L_m, L_n] &= (n - m)L_{n+m} \\ [L_m, Y_p] &= (\rho - m\rho)Y_{m+p} \\ [Y_p, Y_q] &= 0, m, n \in \mathbb{Z}, p, q \in \mathbb{Z} + s. \end{aligned}$$

Schrödinger-Virasoro algebra (Roger-2006, Liu-2016,  
 Li-Su-2008).

$H$ —separable Hilbert space,  $\mathcal{H} = S'(D, H)$ — space of tempered distributions on the unit disk  $D = \{z \in \mathbb{C} \mid |z| \leq 1\}$

$$A_m : \mathcal{H} \mapsto \mathcal{H}, A_m f := e^{-imz}(-m\rho f + i \frac{df}{dz}), z \in D$$

Then

$$[A_m, A_n] = (n - m)A_{n+m}, m, n \in \mathbb{Z},$$

i.e.  $\{A_m, m \in \mathbb{Z}\}$ —representation of Witt algebra. Let  $g \in H$ ,  
 $\{e_p := ge^{-ipz}\}_{p \in \mathbb{Z} + s} \subset \mathcal{H}$ . Then

$$[\bar{D}(A_m), \bar{D}(A_n)] = (n - m)\bar{D}(A_{n+m}), m, n \in \mathbb{Z}.$$

$$[\partial(e_p), \partial(e_q)] = 0, p, q \in \mathbb{Z} + s.$$

$$[\bar{D}(A_m), \partial(e_p)] = \partial(A_m e_p) = (p - m\rho)\partial(e_{m+p}), p \in \mathbb{Z} + s, m \in \mathbb{Z}.$$

$\{\partial(e_p), \bar{D}(A_m), p \in \mathbb{Z} + s, m \in \mathbb{Z}\}$ — a representation of  $\mathcal{L}[s, \rho]$  by linear vector fields.

## Cuntz algebras

$\mathcal{O}_n, n \in \mathbb{N}^*$  – Cuntz algebras (Cuntz–1977) with generators  $\{s_i\}_{i=1,\dots,n}$

$$s_j^* s_j = 1, s_j^* s_k = 0, j \neq k, j, k = 1 \dots, n.$$

$$\sum_{i=1}^n s_i s_i^* = 1 (\leq 1, n = \infty)$$

Algebras of mutually orthogonal isometries of separable Hilbert space.

### Example

$$\mathcal{O}_2 : s_1 e_k := e_{2k}, s_2 e_k := e_{2k+1}$$

## Algebraic analogue of Jordan-Schwinger map

$V$  is separable locally convex TVS in duality with  $W$ ,

biorthogonal system  $\{e_k\}_{k=1}^{\infty} \subset V$ ,  $\{f_k\}_{k=1}^{\infty} \subset W$ .

$$V = \overline{sp\{e_k, k \in \mathbb{N}\}}.$$

$$D(A) := \sum_{\alpha, \beta \in \mathbb{N}} \langle Ae_{\alpha}, f_{\beta} \rangle s_{\alpha} s_{\beta}^*, \quad A \in \mathcal{L}(V, V),$$

$$\partial(h) := \sum_{\alpha \in \mathbb{N}} \langle h, f_{\alpha} \rangle s_{\alpha}^*, \quad h \in V.$$

$$\bar{\partial}(f) := \sum_{\alpha \in I} \langle e_{\alpha}, f \rangle s_{\alpha}, \quad f \in W.$$

# Analogue of J.-S. map: Algebraic homomorphism

## Lemma

$$D(A)D(B) = D(BA),$$

$$\partial(h)D(A) = \partial(Ah),$$

$$D(A)\bar{\partial}(f) = \bar{\partial}(A^*f),$$

$$\partial(h)\bar{\partial}(g) = \langle h, g \rangle$$

$$h \in V, f, g \in W, A, B \in \mathcal{L}(V, V),$$

where the adjoint is taken with respect to the duality  $\langle \cdot, \cdot \rangle$ .

## Invariance of spectrum

$\sigma(A)$ -spectrum of an operator  $A$ ,  $\mathcal{R}(A)$ -resolvent set.

**Corollary (Invariance of the spectrum)**

Assume that the generators  $\{s_i\}$  of  $\mathcal{O}_\infty$  satisfy

$$\sum_i s_i s_i^* = Id.$$

Then  $\sigma(D(A)) \subset \sigma(A)$  and, if  $V$  is a Frechet space,  
 $\sigma(D(A)) = \sigma(A)$ .

# Spectral Theorem

## Corollary (Spectral Theorem)

Assume that  $A \in \mathcal{L}(V, V)$  is symmetric i.e.

$$\langle A\mathbf{e}_\alpha, \mathbf{f}_\beta \rangle = \langle A\mathbf{e}_\beta, \mathbf{f}_\alpha \rangle, \alpha, \beta \in I$$

and  $D(A)$  is selfadjoint (symmetry of  $D(A)$  immediately follows from symmetry of  $A$ ). Then we have

$$\langle Ah, g \rangle = \int_{\sigma(D(A))} \lambda d [\partial(h) E_\lambda \bar{\partial}(g)], \quad h \in V, g \in W,$$

where  $E_\lambda$  denote the spectral projections for the operator  $D(A)$ .

## Representation of algebras

$V = X$  is a topological algebra in duality with  $X^*$ .  $I_a : X \rightarrow X$ ,  
 $I_a(x) = ax$ .  $A, B \in \mathcal{L}(X, X)$ ,  $Ax = I_a x$ ,  $Bx = I_b x$ ,  $a, b \in X$ .

### Corollary

Assume that  $X$  is a topological separable locally convex Hausdorff algebra. Then we have

$$\begin{aligned} D(I_a)D(I_b) &= D(I_{ba}), \\ \partial(h)D(I_a) &= \partial(ah), \\ D(I_a)\bar{\partial}(f) &= \bar{\partial}(I_a^* f), \\ \partial(h)\bar{\partial}(g) &= \langle g, h \rangle \\ a, b, h &\in X, f, g \in X^*. \end{aligned}$$

If  $X$  has identity then  $X \ni a \mapsto D(I_a) \in \mathcal{O}_\infty$  is injective.

# Case of finitedimensional algebra

## Definition (Dutkay-2014)

$Y$  be a topological compact space,  $\mu$  a Borel probability measure on  $Y$ ,  $r : Y \rightarrow Y$  an  $n$ -to-1 Borel measurable map, i.e.  $|r^{-1}(z)| = n$ ,  $\mu$ -a. a.  $z \in Y$ .  $\mu$  is a strongly invariant measure w.r. .  $r$ , i.e.

$$\int f d\mu = \frac{1}{n} \int \sum_{r(\omega)=r(z)} f(\omega) d\mu(z), f \in C(Y)$$

## Definition (Dutkay-2014)

A quadrature mirror filter (QMF) for  $r$  is  $m_0 \in L^\infty(Y, \mu)$

$$\frac{1}{n} \sum_{r(\omega)=z} |m_0(\omega)|^2 = 1, z \in Y.$$

A QMF basis is a set of  $n$  QMF's  $m_0, m_1, \dots, m_{n-1}$

$$\frac{1}{n} \sum_{r(\omega)=z} m_i(\omega) \overline{m_j}(\omega) = 1 \delta_{ij}, i, j \in \{0, 1, \dots, n-1\}, z \in Y.$$

## Case of finitedimensional algebra

**Proposition (Dutkay-2014)**

$\{m_i\}_{i=0}^{i=n-1}$  – QMF basis.

$$S_i(f) = m_i(f \circ r), f \in L^2(Y, d\mu), i = 0, \dots, n-1.$$

Then the operators  $\{S_i\}_{i=0}^{n-1}$  form a representation of  $\mathcal{O}_n$ .

Furthermore,

$$S_i^*(f)(z) = \frac{1}{n} \sum_{r(\omega)=z} \overline{m_i(\omega)} f(\omega), i = 0, \dots, n-1, z \in Y.$$

Consequently,

$$D(I_a)f(z) = \frac{1}{n} \sum_{r(\omega)=r(z)} \left( \sum_{i,j=0}^{n-1} \langle a e_i, f_j \rangle m_i(z) \overline{m_j(\omega)} \right) f(\omega),$$

$$z \in Y, a \in \mathcal{A}, f \in L^2(Y, d\mu)$$

$D(I_a)$  – transfer operator of dynamical system.

# Quantization of finite dimensional systems

$(C^\infty(M), \{\cdot, \cdot\})$  be a Poisson manifold with Poisson bracket  
 $\{\cdot, \cdot\}.$

$$Q(h) := D^*(\{h, \cdot\}), R(h) := D^*(h \cdot), h \in C^\infty(M).$$

## Definition

Quantization  $\widehat{Q} \in \mathcal{L}(C^\infty(M), \mathcal{O}_\infty)$   $\widehat{Q} := R - 2iQ.$

Theorem (Bock,Futorny, N,LMP,2022)

$$\begin{aligned}\widehat{Q}(1) &= Id, \\ [\widehat{Q}(f), \widehat{Q}(g)] &= -2i\widehat{Q}(\{f, g\}), \\ [\widehat{Q}(q_k), \widehat{Q}(q_j)] &= [\widehat{Q}(p_k), \widehat{Q}(p_j)] = 0, \\ [\widehat{Q}(q_k), \widehat{Q}(p_j)] &= -2i\delta_{kj}Id, k, j = 1, \dots, \dim M\end{aligned}$$

Furthermore, if  $\phi : \mathbb{R} \rightarrow \mathbb{R}$  is an analytic function then

$$\Re \widehat{Q}(\phi(f)) = \phi(\Re \widehat{Q}(f))$$

Properties similar to Van Hove prequantization. Generalization to infinite dimensional case in the framework of White noise calculus.

## Summary

Thus we have

- ▶ Class of ergodic system with polynomial rate of convergence
- ▶ Extensions of Jordan-Schwinger map
- ▶ Applications: Quantization, Spectral Theorem...
- ▶ Unsolved problems:
  - ▶ Conservation of  $q$ -bracket.
  - ▶ Jordan-Schwinger map for infinite dimensional algebras such as Yangians, vertex algebras,...
  - ▶ Jordan-Schwinger map as a "linearisation" tool
  - ▶ Connection with dynamical systems

## Literature

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